

# Solutions of nonlinear stochastic differential equations with long-range power-law distributions

Julius Ruseckas, Vygintas Gontis, and Bronsilovas Kaulakys

Institute of Theoretical Physics and Astronomy, Vilnius University, Lithuania

March 17, 2011

Our research is related to the **1/f noise** problem and **long-range** processes.

## 1/f noise

a type of noise whose power spectral density  $S(f)$  behaves like

$$S(f) \sim 1/f^\beta, \quad \beta \text{ is close to } 1$$

Fluctuations of signals exhibiting 1/f behavior of the power spectral density at low frequencies have been observed in a **wide variety** of physical, geophysical, biological, financial, traffic, Internet, astrophysical and other systems.

Our research is related to the **1/f noise** problem and **long-range** processes.

## 1/f noise

a type of noise whose power spectral density  $S(f)$  behaves like

$$S(f) \sim 1/f^\beta, \quad \beta \text{ is close to } 1$$

Fluctuations of signals exhibiting 1/f behavior of the power spectral density at low frequencies have been observed in a **wide variety** of physical, geophysical, biological, financial, traffic, Internet, astrophysical and other systems.

Our research is related to the **1/f noise** problem and **long-range** processes.

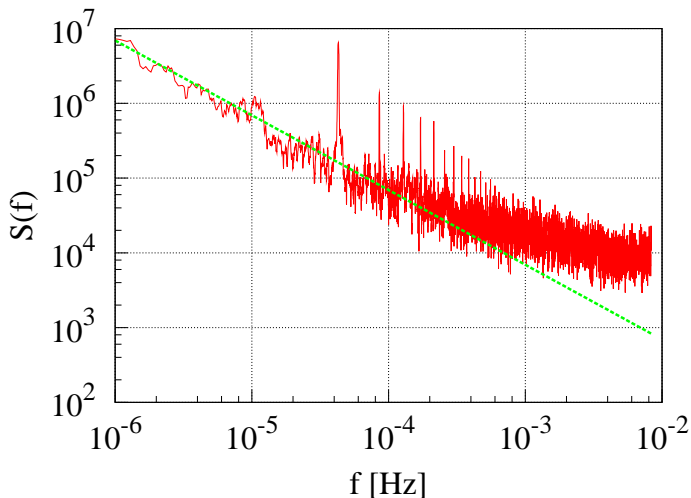
## 1/f noise

a type of noise whose power spectral density  $S(f)$  behaves like

$$S(f) \sim 1/f^\beta, \quad \beta \text{ is close to } 1$$

Fluctuations of signals exhibiting 1/f behavior of the power spectral density at low frequencies have been observed in a **wide variety** of physical, geophysical, biological, financial, traffic, Internet, astrophysical and other systems.

# Example of $1/f$ noise



Power spectral density of trading activity (number of trades per 1 min).  
for ABT stock on NYSE

# Notes about $1/f$ noise

- A **pure**  $1/f$  power spectrum is **physically impossible** because the total power would be infinity.
- We search for a model where the spectrum of signal has  $1/f^\beta$  behavior only in some **intermediate** region of frequencies,  $f_{\min} \ll f \ll f_{\max}$ , whereas for small frequencies  $f \ll f_{\min}$  the spectrum is bounded.
- The behavior of spectrum at frequencies  $f_{\min} \ll f \ll f_{\max}$  is connected with the behavior of the autocorrelation function at times  $1/f_{\max} \ll t \ll 1/f_{\min}$ .

# Notes about $1/f$ noise

- A **pure**  $1/f$  power spectrum is **physically impossible** because the total power would be infinity.
- We search for a model where the spectrum of signal has  $1/f^\beta$  behavior only in some **intermediate** region of frequencies,  $f_{\min} \ll f \ll f_{\max}$ , whereas for small frequencies  $f \ll f_{\min}$  the spectrum is bounded.
- The behavior of spectrum at frequencies  $f_{\min} \ll f \ll f_{\max}$  is connected with the behavior of the autocorrelation function at times  $1/f_{\max} \ll t \ll 1/f_{\min}$ .

# Notes about $1/f$ noise

- A **pure**  $1/f$  power spectrum is **physically impossible** because the total power would be infinity.
- We search for a model where the spectrum of signal has  $1/f^\beta$  behavior only in some **intermediate** region of frequencies,  $f_{\min} \ll f \ll f_{\max}$ , whereas for small frequencies  $f \ll f_{\min}$  the spectrum is bounded.
- The behavior of spectrum at frequencies  $f_{\min} \ll f \ll f_{\max}$  is connected with the behavior of the autocorrelation function at times  $1/f_{\max} \ll t \ll 1/f_{\min}$ .



# Notes about $1/f$ noise

- Often  $1/f$  noise is defined by a long-memory process, characterized by  $S(f) \sim 1/f^\beta$  as  $f \rightarrow 0$ .
- This long-range dependence property is equivalent to similar behavior of autocorrelation function  $C(t)$  as  $t \rightarrow \infty$
- This behavior of the autocorrelation function is **not necessary** for obtaining required form of the power spectrum in a finite interval of the frequencies which does not include zero

# Notes about $1/f$ noise

- Often  $1/f$  noise is defined by a long-memory process, characterized by  $S(f) \sim 1/f^\beta$  as  $f \rightarrow 0$ .
- This long-range dependence property is equivalent to similar behavior of autocorrelation function  $C(t)$  as  $t \rightarrow \infty$
- This behavior of the autocorrelation function is **not necessary** for obtaining required form of the power spectrum in a finite interval of the frequencies which does not include zero

# Notes about $1/f$ noise

- Often  $1/f$  noise is defined by a long-memory process, characterized by  $S(f) \sim 1/f^\beta$  as  $f \rightarrow 0$ .
- This long-range dependence property is equivalent to similar behavior of autocorrelation function  $C(t)$  as  $t \rightarrow \infty$
- This behavior of the autocorrelation function is **not necessary** for obtaining required form of the power spectrum in a finite interval of the frequencies which does not include zero

- $1/f$  noise is intermediate between white noise,  $S(f) \sim 1/f^0$  and Brownian motion  $S(f) \sim 1/f^2$
- In contrast to the Brownian motion generated by the linear stochastic equations, the signals and processes with  $1/f$  spectrum **cannot** be understood and modeled in such a way.

## Goal

to find a simple **nonlinear** stochastic differential equation (SDE) generating signals exhibiting  $1/f$  noise

- $1/f$  noise is intermediate between white noise,  $S(f) \sim 1/f^0$  and Brownian motion  $S(f) \sim 1/f^2$
- In contrast to the Brownian motion generated by the linear stochastic equations, the signals and processes with  $1/f$  spectrum **cannot** be understood and modeled in such a way.

## Goal

to find a simple **nonlinear** stochastic differential equation (SDE) generating signals exhibiting  $1/f$  noise

- $1/f$  noise is intermediate between white noise,  $S(f) \sim 1/f^0$  and Brownian motion  $S(f) \sim 1/f^2$
- In contrast to the Brownian motion generated by the linear stochastic equations, the signals and processes with  $1/f$  spectrum **cannot** be understood and modeled in such a way.

## Goal

to find a simple **nonlinear** stochastic differential equation (SDE) generating signals exhibiting  $1/f$  noise

# Application to the description of trading

- Time series of financial data exhibit highly nontrivial statistical properties. Many of these properties appear to be universal.
- Trading activity, trading volume, and volatility are stochastic variables with the long-range correlation. The autocorrelation of the volatility decays only slowly as a power law.
- Probability distribution functions (PDFs) of return and trading activity have fat tails exhibiting power-law decay.
- Proposed equations can exhibit both power-law PDF and power-law spectrum.

# Application to the description of trading

- Time series of financial data exhibit highly nontrivial statistical properties. Many of these properties appear to be universal.
- Trading activity, trading volume, and volatility are stochastic variables with the long-range correlation. The autocorrelation of the volatility decays only slowly as a power law.
- Probability distribution functions (PDFs) of return and trading activity have fat tails exhibiting power-law decay.
- Proposed equations can exhibit both power-law PDF and power-law spectrum.



# Application to the description of trading

- Time series of financial data exhibit highly nontrivial statistical properties. Many of these properties appear to be universal.
- Trading activity, trading volume, and volatility are stochastic variables with the long-range correlation. The autocorrelation of the volatility decays only slowly as a power law.
- Probability distribution functions (PDFs) of return and trading activity have fat tails exhibiting power-law decay.
- Proposed equations can exhibit both power-law PDF and power-law spectrum.

# Application to the description of trading

- Time series of financial data exhibit highly nontrivial statistical properties. Many of these properties appear to be universal.
- Trading activity, trading volume, and volatility are stochastic variables with the long-range correlation. The autocorrelation of the volatility decays only slowly as a power law.
- Probability distribution functions (PDFs) of return and trading activity have fat tails exhibiting power-law decay.
- Proposed equations can exhibit both power-law PDF and power-law spectrum.

- If  $S(f) \sim f^{-\beta}$  then power spectral density has a scaling property

$$S(af) = a^{-\beta} S(f)$$

- Wiener-Khintchine theorem

$$C(t) = \int_{-\infty}^{+\infty} S(f) \cos(2\pi ft) df$$

- Autocorrelation function  $C(t)$  has scaling property

$$C(at) \sim a^{\beta-1} C(t)$$

- If  $S(f) \sim f^{-\beta}$  then power spectral density has a scaling property

$$S(af) = a^{-\beta} S(f)$$

- Wiener-Khintchine theorem

$$C(t) = \int_{-\infty}^{+\infty} S(f) \cos(2\pi ft) df$$

- Autocorrelation function  $C(t)$  has scaling property

$$C(at) \sim a^{\beta-1} C(t)$$

- If  $S(f) \sim f^{-\beta}$  then power spectral density has a scaling property

$$S(af) = a^{-\beta} S(f)$$

- Wiener-Khintchine theorem

$$C(t) = \int_{-\infty}^{+\infty} S(f) \cos(2\pi ft) df$$

- Autocorrelation function  $C(t)$  has scaling property

$$C(at) \sim a^{\beta-1} C(t)$$

- Autocorrelation function can be written as

$$C(t) = \int dx \int dx' xx' P_0(x) P_x(x', t|x, 0)$$

- $P_0(x)$  is the steady state PDF
- $P_x(x', t|x, 0)$  is the transition probability
- The transition probability can be obtained from the solution of the **Fokker-Planck equation** with the initial condition  $P_x(x', 0|x, 0) = \delta(x' - x)$ .

- Autocorrelation function can be written as

$$C(t) = \int dx \int dx' xx' P_0(x) P_x(x', t|x, 0)$$

- $P_0(x)$  is the steady state PDF
- $P_x(x', t|x, 0)$  is the transition probability
- The transition probability can be obtained from the solution of the **Fokker-Planck equation** with the initial condition  $P_x(x', 0|x, 0) = \delta(x' - x)$ .

- Autocorrelation function can be written as

$$C(t) = \int dx \int dx' xx' P_0(x) P_x(x', t|x, 0)$$

- $P_0(x)$  is the steady state PDF
- $P_x(x', t|x, 0)$  is the transition probability
- The transition probability can be obtained from the solution of the **Fokker-Planck equation** with the initial condition  $P_x(x', 0|x, 0) = \delta(x' - x)$ .



- Autocorrelation function can be written as

$$C(t) = \int dx \int dx' xx' P_0(x) P_x(x', t|x, 0)$$

- $P_0(x)$  is the steady state PDF
- $P_x(x', t|x, 0)$  is the transition probability
- The transition probability can be obtained from the solution of the **Fokker-Planck equation** with the initial condition  $P_x(x', 0|x, 0) = \delta(x' - x)$ .

- Let us assume that

- Steady state PDF has power-law form

$$P_0(x) \sim x^{-\nu}$$

- Transition probability has a scaling property

$$P(ax', t|ax, 0) = a^{-1} P(x', a^{2(\eta-1)}t|x, 0)$$

- Then the autocorrelation function will have the required scaling with

$$\beta = 1 + \frac{\nu - 3}{2(\eta - 1)}$$

# Justification of SDE

- Let us assume that
  - Steady state PDF has power-law form

$$P_0(x) \sim x^{-\nu}$$

- Transition probability has a scaling property

$$P(ax', t|ax, 0) = a^{-1} P(x', a^{2(\eta-1)}t|x, 0)$$

- Then the autocorrelation function will have the required scaling with

$$\beta = 1 + \frac{\nu - 3}{2(\eta - 1)}$$

# Justification of SDE

- Let us assume that
  - Steady state PDF has power-law form

$$P_0(x) \sim x^{-\nu}$$

- Transition probability has a scaling property

$$P(ax', t|ax, 0) = a^{-1} P(x', a^{2(\eta-1)}t|x, 0)$$

- Then the autocorrelation function will have the required scaling with

$$\beta = 1 + \frac{\nu - 3}{2(\eta - 1)}$$

- Let us assume that
  - Steady state PDF has power-law form

$$P_0(x) \sim x^{-\nu}$$

- Transition probability has a scaling property

$$P(ax', t|ax, 0) = a^{-1} P(x', a^{2(\eta-1)}t|x, 0)$$

- Then the autocorrelation function will have the required scaling with

$$\beta = 1 + \frac{\nu - 3}{2(\eta - 1)}$$

# Justification of SDE

To get the required scaling of transition probability:

- SDE will contain only powers of  $x$
- The diffusion coefficient will be of the form  $x^{2\eta}$
- The drift term is fixed by the requirement that the steady-state PDF should be  $x^{-\nu}$

## Proposed SDE

$$dx = \sigma^2(\eta - \nu/2)x^{2\eta-1}dt + \sigma x^\eta dW_t$$

B. Kaulakys and J. Ruseckas, Phys. Rev. E **70**, 020101(R) (2004).

B. Kaulakys and J. Ruseckas, V. Gontis, and M. Alaburda, Physica A **365**, 217 (2006).

# Justification of SDE

To get the required scaling of transition probability:

- SDE will contain only powers of  $x$
- The diffusion coefficient will be of the form  $x^{2\eta}$
- The drift term is fixed by the requirement that the steady-state PDF should be  $x^{-\nu}$

## Proposed SDE

$$dx = \sigma^2(\eta - \nu/2)x^{2\eta-1}dt + \sigma x^\eta dW_t$$

B. Kaulakys and J. Ruseckas, Phys. Rev. E **70**, 020101(R) (2004).

B. Kaulakys and J. Ruseckas, V. Gontis, and M. Alaburda, Physica A **365**, 217 (2006).

# Justification of SDE

To get the required scaling of transition probability:

- SDE will contain only powers of  $x$
- The diffusion coefficient will be of the form  $x^{2\eta}$
- The drift term is fixed by the requirement that the steady-state PDF should be  $x^{-\nu}$

## Proposed SDE

$$dx = \sigma^2(\eta - \nu/2)x^{2\eta-1}dt + \sigma x^\eta dW_t$$

B. Kaulakys and J. Ruseckas, Phys. Rev. E **70**, 020101(R) (2004).

B. Kaulakys and J. Ruseckas, V. Gontis, and M. Alaburda, Physica A **365**, 217 (2006).



# Justification of SDE

To get the required scaling of transition probability:

- SDE will contain only powers of  $x$
- The diffusion coefficient will be of the form  $x^{2\eta}$
- The drift term is fixed by the requirement that the steady-state PDF should be  $x^{-\nu}$

## Proposed SDE

$$dx = \sigma^2(\eta - \nu/2)x^{2\eta-1}dt + \sigma x^\eta dW_t$$

B. Kaulakys and J. Ruseckas, Phys. Rev. E **70**, 020101(R) (2004).

B. Kaulakys and J. Ruseckas, V. Gontis, and M. Alaburda, Physica A **365**, 217 (2006).

# Justification of SDE

To get the required scaling of transition probability:

- SDE will contain only powers of  $x$
- The diffusion coefficient will be of the form  $x^{2\eta}$
- The drift term is fixed by the requirement that the steady-state PDF should be  $x^{-\nu}$

## Proposed SDE

$$dx = \sigma^2(\eta - \nu/2)x^{2\eta-1}dt + \sigma x^\eta dW_t$$

B. Kaulakys and J. Ruseckas, Phys. Rev. E **70**, 020101(R) (2004).

B. Kaulakys and J. Ruseckas, V. Gontis, and M. Alaburda, Physica A **365**, 217 (2006).

# Restriction of diffusion

- Because of the divergence of the power-law distribution and the requirement of the stationarity of the process, the SDE should be analyzed together with the appropriate **restrictions** of the diffusion in some finite interval.
- When diffusion is restricted, **scaling properties are only approximate**, but  $1/f$  spectrum remains in a wide interval of frequencies.

# Restriction of diffusion

- Because of the divergence of the power-law distribution and the requirement of the stationarity of the process, the SDE should be analyzed together with the appropriate **restrictions** of the diffusion in some finite interval.
- When diffusion is restricted, **scaling properties are only approximate**, but  $1/f$  spectrum remains in a wide interval of frequencies.

# Restriction of diffusion

Possible forms of restriction:

- Reflective boundary conditions at  $x = x_{\min}$  and  $x = x_{\max}$
- Exponential restriction of the diffusion

$$dx = \sigma^2 \left( \eta - \frac{\nu}{2} + \frac{m}{2} \left( \frac{x_{\min}}{x} \right)^m - \frac{m}{2} \left( \frac{x}{x_{\max}} \right)^m \right) x^{2\eta-1} dt + \sigma x^\eta dW_t$$

Steady state PDF:

$$P_0(x) \sim x^{-\nu} \exp \left( - \left( \frac{x_{\min}}{x} \right)^m - \left( \frac{x}{x_{\max}} \right)^m \right)$$

# Restriction of diffusion

Possible forms of restriction:

- Reflective boundary conditions at  $x = x_{\min}$  and  $x = x_{\max}$
- Exponential restriction of the diffusion

$$dx = \sigma^2 \left( \eta - \frac{\nu}{2} + \frac{m}{2} \left( \frac{x_{\min}}{x} \right)^m - \frac{m}{2} \left( \frac{x}{x_{\max}} \right)^m \right) x^{2\eta-1} dt + \sigma x^\eta dW_t$$

Steady state PDF:

$$P_0(x) \sim x^{-\nu} \exp \left( - \left( \frac{x_{\min}}{x} \right)^m - \left( \frac{x}{x_{\max}} \right)^m \right)$$

Possible forms of restriction:

- Reflective boundary conditions at  $x = x_{\min}$  and  $x = x_{\max}$
- Exponential restriction of the diffusion

$$dx = \sigma^2 \left( \eta - \frac{\nu}{2} + \frac{m}{2} \left( \frac{x_{\min}}{x} \right)^m - \frac{m}{2} \left( \frac{x}{x_{\max}} \right)^m \right) x^{2\eta-1} dt + \sigma x^\eta dW_t$$

Steady state PDF:

$$P_0(x) \sim x^{-\nu} \exp \left( - \left( \frac{x_{\min}}{x} \right)^m - \left( \frac{x}{x_{\max}} \right)^m \right)$$

# Restriction of diffusion

- $q$ -exponential steady-state PDF

$$dx = \sigma^2(\eta - \nu/2)(x + x_0)^{2\eta-1}dt + \sigma(x + x_0)^\eta dW_t$$

$$P_0(x) \sim \exp_{1+1/\nu}(-\nu x/x_0)$$

Reflective boundary condition at  $x = 0$

B. Kaulakys and M. Alaburda, J. Stat. Mech. **2009**, P02051 (2009).

- $q$ -Gaussian steady-state PDF

$$dx = \sigma^2(\eta - \nu/2)(x^2 + x_0^2)^{\eta-1}xdt + \sigma(x^2 + x_0^2)^{\eta/2}dW_t$$

$$P_0(x) \sim \exp_{1+2/\nu}(-\nu x^2/2x_0^2)$$

B. Kaulakys, M. Alaburda, and V. Gontis, AIP Conf. Proc. **1129**, 13 (2009).

V. Gontis, B. Kaulakys, and J. Ruseckas, AIP Conf. Proc. **1129**, 563 (2009).

V. Gontis, J. Ruseckas, and A. Kononovičius, Physica A, **389**, 100 (2010).

$q$ -exponential function:  $\exp_q(x) \equiv (1 + (1 - q)x)^{1/(1-q)}$



# Restriction of diffusion

- $q$ -exponential steady-state PDF

$$dx = \sigma^2(\eta - \nu/2)(x + x_0)^{2\eta-1}dt + \sigma(x + x_0)^\eta dW_t$$

$$P_0(x) \sim \exp_{1+1/\nu}(-\nu x/x_0)$$

Reflective boundary condition at  $x = 0$

B. Kaulakys and M. Alaburda, J. Stat. Mech. **2009**, P02051 (2009).

- $q$ -Gaussian steady-state PDF

$$dx = \sigma^2(\eta - \nu/2)(x^2 + x_0^2)^{\eta-1}xdt + \sigma(x^2 + x_0^2)^{\eta/2}dW_t$$

$$P_0(x) \sim \exp_{1+2/\nu}(-\nu x^2/2x_0^2)$$

B. Kaulakys, M. Alaburda, and V. Gontis, AIP Conf. Proc. **1129**, 13 (2009).

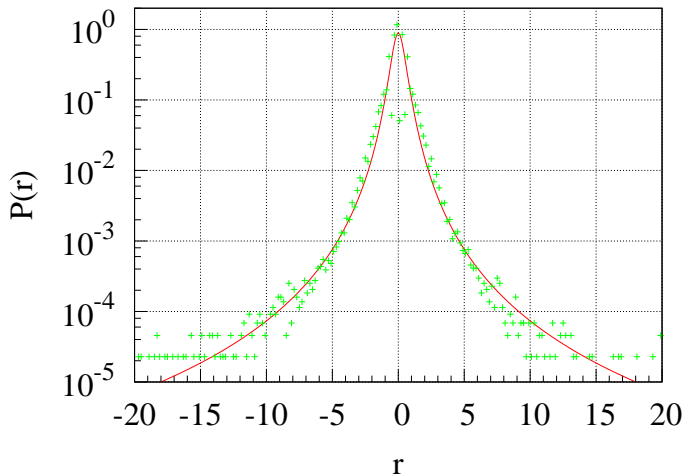
V. Gontis, B. Kaulakys, and J. Ruseckas, AIP Conf. Proc. **1129**, 563 (2009).

V. Gontis, J. Ruseckas, and A. Kononovičius, Physica A, **389**, 100 (2010).

$q$ -exponential function:  $\exp_q(x) \equiv (1 + (1 - q)x)^{1/(1-q)}$

# Trading return

The distribution of normalized return  $r$  per 1 min is close to  $q$ -Gaussian.



Normalized trading return per 1 min for ABT stock on NYSE

# Connection with other equations

For some choices of parameters our SDE takes the form of well-known SDE's considered in econophysics and finance.

- $\eta = 0$  and  $\sigma = 1$  corresponds to the **Bessel process**

$$dx = \frac{\delta - 1}{2} \frac{1}{x} dt + dW_t$$

of dimension  $\delta = 1 - \nu$

- $\eta = 1/2$ ,  $\sigma = 2$  corresponds to the **squared Bessel process**

$$dx = \delta dt + 2\sqrt{x} dW_t$$

of dimension  $\delta = 2(1 - \nu)$

# Connection with other equations

For some choices of parameters our SDE takes the form of well-known SDE's considered in econophysics and finance.

- $\eta = 0$  and  $\sigma = 1$  corresponds to the **Bessel process**

$$dx = \frac{\delta - 1}{2} \frac{1}{x} dt + dW_t$$

of dimension  $\delta = 1 - \nu$

- $\eta = 1/2$ ,  $\sigma = 2$  corresponds to the **squared Bessel process**

$$dx = \delta dt + 2\sqrt{x} dW_t$$

of dimension  $\delta = 2(1 - \nu)$

# Connection with other equations

For some choices of parameters our SDE takes the form of well-known SDE's considered in econophysics and finance.

- $\eta = 0$  and  $\sigma = 1$  corresponds to the **Bessel process**

$$dx = \frac{\delta - 1}{2} \frac{1}{x} dt + dW_t$$

of dimension  $\delta = 1 - \nu$

- $\eta = 1/2$ ,  $\sigma = 2$  corresponds to the **squared Bessel process**

$$dx = \delta dt + 2\sqrt{x} dW_t$$

of dimension  $\delta = 2(1 - \nu)$

# Connection with other equations

- SDE with exponential restriction with  $\eta = 1/2$ ,  $x_{\min} = 0$  and  $m = 1$  gives **Cox-Ingersoll-Ross (CIR) process**

$$dx = k(\theta - x)dt + \sigma\sqrt{x}dW_t$$

where  $k = \sigma^2/2x_{\max}$ ,  $\theta = x_{\max}(1 - \nu)$

- When  $\nu = 2\eta$ ,  $x_{\max} = \infty$  and  $m = 2\eta - 2$  then we get the **Constant Elasticity of Variance (CEV) process**

$$dx = \mu x dt + \sigma x^\eta dW_t$$

where  $\mu = \sigma^2(\eta - 1)x_{\min}^{2(\eta-1)}$

- SDE with exponential restriction with  $\eta = 1/2$ ,  $x_{\min} = 0$  and  $m = 1$  gives **Cox-Ingersoll-Ross (CIR) process**

$$dx = k(\theta - x)dt + \sigma\sqrt{x}dW_t$$

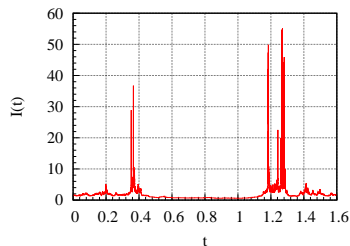
where  $k = \sigma^2/2x_{\max}$ ,  $\theta = x_{\max}(1 - \nu)$

- When  $\nu = 2\eta$ ,  $x_{\max} = \infty$  and  $m = 2\eta - 2$  then we get the **Constant Elasticity of Variance (CEV) process**

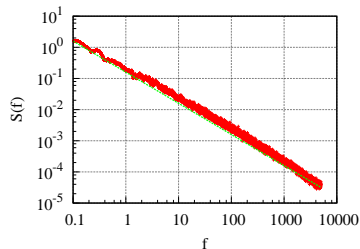
$$dx = \mu x dt + \sigma x^\eta dW_t$$

where  $\mu = \sigma^2(\eta - 1)x_{\min}^{2(\eta-1)}$

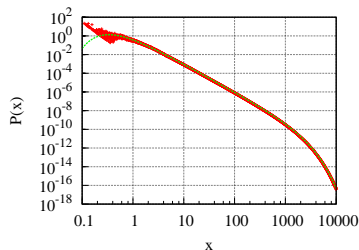
# Numerical simulation



Typical signal



Power spectral density



Distribution of  $x$

Used parameters:  $\nu = 3$ ,  $\eta = 5/2$ ,  
 $x_{\min} = 1.0$ ,  $x_{\max} = 10^3$ .  
 $1/f$  spectrum.



# Analytically solvable case

A CEV process:

$$dx = \mu x dt + \sigma x^{\frac{3}{2}} dW_t$$

where  $\mu = \sigma^2 x_{\min}/2$ ,  $\eta = 3/2$ ,  $\nu = 3$  and  $x_{\max} = \infty$

Transition probability is

$$P_x(x', t|x, 0) = \frac{x_{\min}}{(1 - e^{-\mu t})} \sqrt{\frac{x}{x'^5}} \exp\left(\frac{1}{2}\mu t - \frac{x_{\min}}{(1 - e^{-\mu t})} \left(\frac{1}{x'} + \frac{1}{x} e^{-\mu t}\right)\right) \\ \times I_1\left(\frac{x_{\min}}{\sinh(\frac{1}{2}\mu t)} \frac{1}{\sqrt{xx'}}\right)$$

The steady-state probability distribution has the form

$$P_0(x) = x_{\min}^2 x^{-3} \exp(-x_{\min}/x)$$

# Analytically solvable case

A CEV process:

$$dx = \mu x dt + \sigma x^{\frac{3}{2}} dW_t$$

where  $\mu = \sigma^2 x_{\min}/2$ ,  $\eta = 3/2$ ,  $\nu = 3$  and  $x_{\max} = \infty$

Transition probability is

$$P_x(x', t|x, 0) = \frac{x_{\min}}{(1 - e^{-\mu t})} \sqrt{\frac{x}{x'^5}} \exp\left(\frac{1}{2}\mu t - \frac{x_{\min}}{(1 - e^{-\mu t})} \left(\frac{1}{x'} + \frac{1}{x} e^{-\mu t}\right)\right) \\ \times I_1\left(\frac{x_{\min}}{\sinh(\frac{1}{2}\mu t)} \frac{1}{\sqrt{xx'}}\right)$$

The steady-state probability distribution has the form

$$P_0(x) = x_{\min}^2 x^{-3} \exp(-x_{\min}/x)$$

# Analytically solvable case

The autocorrelation function

$$C(t) = -x_{\min}^2 e^{\mu t} \ln(1 - e^{-\mu t})$$

When  $\mu t \ll 1$  we get  $C(t) \approx -x_{\min}^2 \ln(\mu t)$

The power spectral density

$$S(f) = -4x_{\min}^2 \operatorname{Re} \left[ \frac{\gamma + \psi(i\omega/\mu)}{\mu - i\omega} \right]$$

where  $\gamma$  is the Euler's constant and  $\psi(\cdot)$  is the digamma function.

When  $\omega \gg \mu$  then the power spectral density is  $S(f) \approx x_{\min}^2/f$

# Analytically solvable case

The autocorrelation function

$$C(t) = -x_{\min}^2 e^{\mu t} \ln(1 - e^{-\mu t})$$

When  $\mu t \ll 1$  we get  $C(t) \approx -x_{\min}^2 \ln(\mu t)$

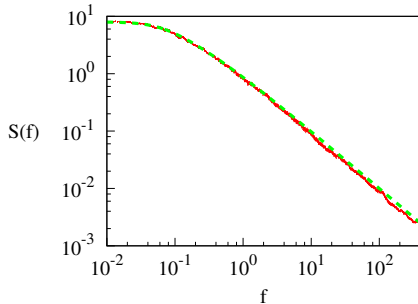
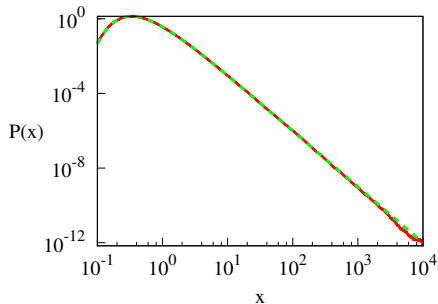
The power spectral density

$$S(f) = -4x_{\min}^2 \operatorname{Re} \left[ \frac{\gamma + \psi(i\omega/\mu)}{\mu - i\omega} \right]$$

where  $\gamma$  is the Euler's constant and  $\psi(\cdot)$  is the digamma function.

When  $\omega \gg \mu$  then the power spectral density is  $S(f) \approx x_{\min}^2/f$

# Analytically solvable case



Probability distribution function  $P_0(x)$  and power spectral density  $S(f)$

# 1/f noise and eigenvalues of the F-P equation

- Solutions of the Fokker-Planck equation having the form  $P(x, t) = P_\lambda(x)e^{-\lambda t}$  determine eigenfunctions  $P_\lambda(x)$  and eigenvalues  $\lambda$
- The power spectral density

$$S(f) = 4 \sum_{\lambda} \frac{\lambda}{\lambda^2 + \omega^2} X_{\lambda}^2, \quad X_{\lambda} = \int_{x_{\min}}^{x_{\max}} x P_{\lambda}(x) dx$$

- The shape of the power spectral density depends on the behavior of the eigenfunctions and the eigenvalues
- Expression for the power spectral density resembles the models of 1/f noise using the sum of the Lorentzian spectra. The Lorentzians can arise from the single nonlinear stochastic differential equation

# 1/f noise and eigenvalues of the F-P equation

- Solutions of the Fokker-Planck equation having the form  $P(x, t) = P_\lambda(x)e^{-\lambda t}$  determine eigenfunctions  $P_\lambda(x)$  and eigenvalues  $\lambda$
- The power spectral density

$$S(f) = 4 \sum_{\lambda} \frac{\lambda}{\lambda^2 + \omega^2} X_{\lambda}^2, \quad X_{\lambda} = \int_{x_{\min}}^{x_{\max}} x P_{\lambda}(x) dx$$

- The shape of the power spectral density depends on the behavior of the eigenfunctions and the eigenvalues
- Expression for the power spectral density resembles the models of 1/f noise using the sum of the Lorentzian spectra. The Lorentzians can arise from the single nonlinear stochastic differential equation

# 1/f noise and eigenvalues of the F-P equation

- Solutions of the Fokker-Planck equation having the form  $P(x, t) = P_\lambda(x)e^{-\lambda t}$  determine eigenfunctions  $P_\lambda(x)$  and eigenvalues  $\lambda$
- The power spectral density

$$S(f) = 4 \sum_{\lambda} \frac{\lambda}{\lambda^2 + \omega^2} X_{\lambda}^2, \quad X_{\lambda} = \int_{x_{\min}}^{x_{\max}} x P_{\lambda}(x) dx$$

- The shape of the power spectral density depends on the behavior of the eigenfunctions and the eigenvalues
- Expression for the power spectral density resembles the models of 1/f noise using the sum of the Lorentzian spectra. The Lorentzians can arise from the single nonlinear stochastic differential equation



# 1/f noise and eigenvalues of the F-P equation

- Solutions of the Fokker-Planck equation having the form  $P(x, t) = P_\lambda(x)e^{-\lambda t}$  determine eigenfunctions  $P_\lambda(x)$  and eigenvalues  $\lambda$
- The power spectral density

$$S(f) = 4 \sum_{\lambda} \frac{\lambda}{\lambda^2 + \omega^2} X_{\lambda}^2, \quad X_{\lambda} = \int_{x_{\min}}^{x_{\max}} x P_{\lambda}(x) dx$$

- The shape of the power spectral density depends on the behavior of the eigenfunctions and the eigenvalues
- Expression for the power spectral density resembles the models of 1/f noise using the sum of the Lorentzian spectra. The Lorentzians can arise from the single nonlinear stochastic differential equation

# 1/f noise and eigenvalues of the F-P equation

$$S(f) \approx 4 \int \frac{\lambda}{\lambda^2 + \omega^2} X_\lambda^2 D(\lambda) d\lambda \sim \int_{\lambda_{\min}}^{\lambda_{\max}} \frac{1}{\lambda^{\beta-1}} \frac{1}{\lambda^2 + \omega^2} d\lambda$$

The largest contribution make the terms corresponding to the eigenvalues  $\lambda$  obeying the condition  $\lambda_{\min} \ll \lambda \ll \lambda_{\max}$ , where

$$\begin{aligned} \lambda_{\min} &= \sigma^2 X_{\min}^{2(\eta-1)}, & \lambda_{\max} &= \sigma^2 X_{\max}^{2(\eta-1)}, & \eta > 1 \\ \lambda_{\min} &= \sigma^2 X_{\max}^{2(\eta-1)}, & \lambda_{\max} &= \sigma^2 X_{\min}^{2(\eta-1)}, & \eta < 1 \end{aligned}$$

When  $\lambda_{\min} \ll \omega \ll \lambda_{\max}$  then the leading term in the expansion in the power series of  $\omega$  is

$$S(f) \sim \omega^{-\beta}, \quad \beta < 2$$

# 1/f noise and eigenvalues of the F-P equation

$$S(f) \approx 4 \int \frac{\lambda}{\lambda^2 + \omega^2} X_\lambda^2 D(\lambda) d\lambda \sim \int_{\lambda_{\min}}^{\lambda_{\max}} \frac{1}{\lambda^{\beta-1}} \frac{1}{\lambda^2 + \omega^2} d\lambda$$

The largest contribution make the terms corresponding to the eigenvalues  $\lambda$  obeying the condition  $\lambda_{\min} \ll \lambda \ll \lambda_{\max}$ , where

$$\begin{aligned} \lambda_{\min} &= \sigma^2 X_{\min}^{2(\eta-1)}, & \lambda_{\max} &= \sigma^2 X_{\max}^{2(\eta-1)}, & \eta > 1 \\ \lambda_{\min} &= \sigma^2 X_{\max}^{2(\eta-1)}, & \lambda_{\max} &= \sigma^2 X_{\min}^{2(\eta-1)}, & \eta < 1 \end{aligned}$$

When  $\lambda_{\min} \ll \omega \ll \lambda_{\max}$  then the leading term in the expansion in the power series of  $\omega$  is

$$S(f) \sim \omega^{-\beta}, \quad \beta < 2$$

# $1/f$ noise and eigenvalues of the F-P equation

- The shape of the power spectral density depends on the behavior of the eigenfunctions and the eigenvalues in terms of the function  $X_\lambda^2 D(\lambda)$ .
- One obtains  $1/f^\beta$  behavior of the power spectral density when function  $X_\lambda^2 D(\lambda)$  is proportional to  $\lambda^{-\beta}$  for a wide range of eigenvalues  $\lambda$

J. Ruseckas and B. Kaulakys, Phys. Rev. E **81**, 031105 (2010).

# $1/f$ noise and eigenvalues of the F-P equation

- The shape of the power spectral density depends on the behavior of the eigenfunctions and the eigenvalues in terms of the function  $X_\lambda^2 D(\lambda)$ .
- One obtains  $1/f^\beta$  behavior of the power spectral density when function  $X_\lambda^2 D(\lambda)$  is proportional to  $\lambda^{-\beta}$  for a wide range of eigenvalues  $\lambda$

J. Ruseckas and B. Kaulakys, Phys. Rev. E **81**, 031105 (2010).

# Summary

- We obtain a class of nonlinear SDEs, giving the power-law behavior of the power spectral density in any desirably wide range of frequencies
- and power-law steady state distribution of the signal intensity.
- The equations, as special cases, contain the well-known SDEs in economics and finance.
- One of the reasons for the appearance of the  $1/f$  spectrum is the scaling property of the SDE.
- The power spectral density may be represented as a sum of the Lorentzian spectra.

# Summary

- We obtain a class of nonlinear SDEs, giving the power-law behavior of the power spectral density in any desirably wide range of frequencies
- and power-law steady state distribution of the signal intensity.
- The equations, as special cases, contain the well-known SDEs in economics and finance.
- One of the reasons for the appearance of the  $1/f$  spectrum is the scaling property of the SDE.
- The power spectral density may be represented as a sum of the Lorentzian spectra.

# Summary

- We obtain a class of nonlinear SDEs, giving the power-law behavior of the power spectral density in any desirably wide range of frequencies
- and power-law steady state distribution of the signal intensity.
- The equations, as special cases, contain the well-known SDEs in economics and finance.
- One of the reasons for the appearance of the  $1/f$  spectrum is the scaling property of the SDE.
- The power spectral density may be represented as a sum of the Lorentzian spectra.



# Summary

- We obtain a class of nonlinear SDEs, giving the power-law behavior of the power spectral density in any desirably wide range of frequencies
- and power-law steady state distribution of the signal intensity.
- The equations, as special cases, contain the well-known SDEs in economics and finance.
- One of the reasons for the appearance of the  $1/f$  spectrum is the scaling property of the SDE.
- The power spectral density may be represented as a sum of the Lorentzian spectra.

- We obtain a class of nonlinear SDEs, giving the power-law behavior of the power spectral density in any desirably wide range of frequencies
- and power-law steady state distribution of the signal intensity.
- The equations, as special cases, contain the well-known SDEs in economics and finance.
- One of the reasons for the appearance of the  $1/f$  spectrum is the scaling property of the SDE.
- The power spectral density may be represented as a sum of the Lorentzian spectra.

# Thank you for your attention!